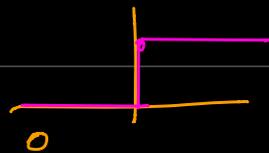


Osservazione sulla soluzione del problema di Cauchy

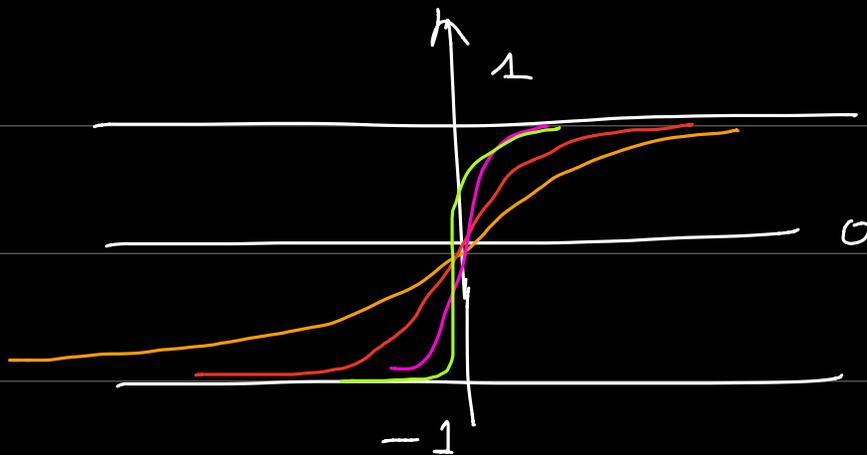
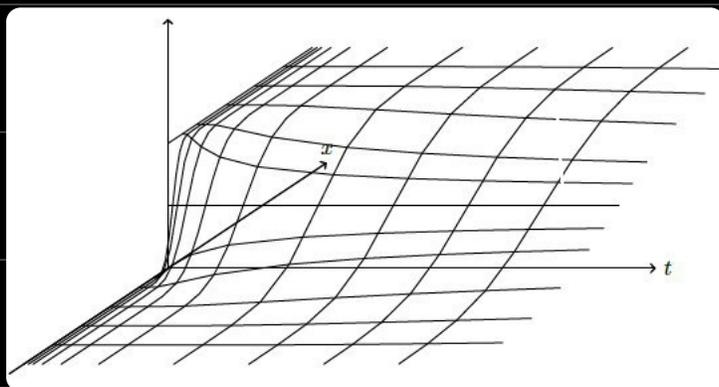
$$\begin{cases} u_t = \Delta u \\ u(0, x) = g(x) = \begin{cases} 1 & x \geq 0 \\ 0 & x < 0 \end{cases} \end{cases}$$



$$u(t, x) = \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} e^{-\frac{|x-y|^2}{4t}} g(y) dy = \frac{1}{\sqrt{4\pi t}} \int_0^{+\infty} e^{-\frac{|x-y|^2}{4t}} dy$$

$$\frac{x-y}{\sqrt{4t}} = s \Rightarrow \frac{1}{\sqrt{\pi}} \int_0^{\frac{x}{\sqrt{4t}}} e^{-s^2} ds = \frac{1}{2} \left(\text{Erf}\left(\frac{x}{\sqrt{4t}}\right) + 1 \right)$$

$$\text{Erf}(x) := \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$$



ESERCIZI

① Dimostrare che se u risolve

$$\begin{cases} u_t = \Delta u & \mathbb{R}^n \times \mathbb{R} \\ u(0, x) = u_0(x) & \mathbb{R} \end{cases}$$

Allora $v(t, x) = \int_0^x u(t, s) ds$ e $w(t, x) = u_x(t, x)$ soddisfano

la stessa equazione con dati iniziali spettivamente integrati e derivati. Dobbiamo supporre che questi ($e u_0$) siano limitati, in modo da poter utilizzare anche per questi la formula di rappresentazione.

Sviluppo

$$u(t, x) = \int_{\mathbb{R}} \Phi(t, x-y) u_0(y) dy \Rightarrow \partial_x u = \int_{\mathbb{R}} \partial_x \Phi(t, x-y) u_0(y) dy$$

e consideriamo

$$w(t, x) = \int_{\mathbb{R}} \Phi(t, x-y) u_0'(y) dy$$

Integriamo per parti per $t > 0$ e x fissati:

$$\begin{aligned} \int_{\mathbb{R}} \Phi(t, x-y) u_0'(y) dy &= \lim_{R \rightarrow +\infty} \int_{-R}^R \Phi(t, x-y) u_0'(y) dy = \\ &= \lim_{R \rightarrow +\infty} \left\{ \Phi(t, x-y) u_0(y) \Big|_{-R}^R + \int_{-R}^R \Phi_x(t, x-y) u_0(y) dy \right\} \\ &= \int_{\mathbb{R}} \Phi_x(t, x-y) u_0(y) dy = \partial_x u(t, x) \end{aligned}$$

OSSERVAZIONE IMPORTANTE: abbiamo dimostrato che

$$(\Phi * f)' = \Phi' * f = \Phi * f'$$

è una proprietà fondamentale dei prodotti di convoluzione

Per l'integrale ragioniamo in modo del tutto analogo

- Abbiamo già visto che la soluzione di:

$$\begin{cases} u_t = \Delta u \\ u(0, x) = e^{-x^2} = g_0(x) \end{cases}$$

$$è \quad u(t, x) = \frac{1}{\sqrt{1+4t}} e^{-\frac{x^2}{1+4t}}$$



$$\text{Erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt \quad \text{Quindi la soluzione di}$$

$$\begin{cases} v_t = \Delta v \\ v(0, x) = \text{Erf}(x) = g_1(x) \end{cases}$$

$$\text{sarà} \quad v_1(t, x) = \frac{2}{\sqrt{\pi}} \int_0^x \frac{1}{\sqrt{1+4t}} e^{-\frac{s^2}{1+4t}} ds = \frac{2}{\sqrt{\pi}} \text{Erf}\left(\frac{x}{\sqrt{1+4t}}\right)$$

$$- \quad g_1(x) = x e^{-x^2} = -\frac{1}{2} \frac{d}{dx} e^{-x^2}$$

⇒ la soluzione corrispondente sarà

$$v_2(t, x) = -\frac{1}{2} \frac{d}{dx} \left(\frac{1}{\sqrt{1+4t}} e^{-\frac{x^2}{1+4t}} \right) = \frac{1}{(1+4t)^{3/2}} x e^{-\frac{x^2}{1+4t}}$$

se deriviamo ancora una volta in x otteniamo

$$w(t, x) = -\frac{2}{(1+4t)^{5/2}} x^2 e^{-\frac{x^2}{1+4t}} + \frac{1}{(1+4t)^{3/2}} e^{-\frac{x^2}{1+4t}}$$

che ha come dato iniziale

$$-2x^2 e^{-x^2} + e^{-x^2} = g_1 + g_0$$

Quindi, la soluzione

$$u_3(t, x) = \frac{1}{(1+4t)^{5/2}} x^2 e^{-\frac{x^2}{1+4t}} - \frac{1}{2(1+4t)^{3/2}} e^{-\frac{x^2}{1+4t}} + \frac{1}{2(1+4t)^{1/2}} e^{-\frac{x^2}{1+4t}}$$

(2) Dimostrare, utilizzando la formula di rappresentazione

$$\text{che } \int_{\mathbb{R}} x u(t, x) dx = \int_{\mathbb{R}} x u_0(x) dx$$

se $x u_0(x) \in L^1$

Scriviamo:

$$\int_{\mathbb{R}} x u(t, x) dx = \int_{\mathbb{R}} x \int_{\mathbb{R}} \Phi(t, x-y) u_0(y) dy dx$$

$$= \iint_{\mathbb{R} \times \mathbb{R}} x \Phi(t, x-y) u_0(y) dy dx$$

$$= \iint_{\mathbb{R} \times \mathbb{R}} [(x-y) \Phi(t, x-y) u_0(y) + \Phi(t, x-y) y u_0(y)] dy dx$$

$$= \iint_{\mathbb{R} \times \mathbb{R}} (x-y) \Phi(t, x-y) u_0(y) dx dy + \iint_{\mathbb{R} \times \mathbb{R}} \Phi(t, x-y) u_0(y) dy dx$$

$$= \int_{\mathbb{R}} y u_0(y) dy \quad \int_{\mathbb{R}} \Phi(t, x) dx = 1 \quad \int_{\mathbb{R}} x \Phi(t, x) dx = 0$$

↓ per sim x

possiamo applicare il teorema di Fubini-Tonelli

Possiamo ragionare in modo analogo per

$$\int_{\mathbb{R}} x^2 u(t, x) dx = \int_{\mathbb{R} \times \mathbb{R}} x^2 \Phi(t, x-y) u_0(y) dx$$

$$x^2 = (x-y)^2 + 2y(x-y) + y^2$$

$$= \iint_{\mathbb{R} \times \mathbb{R}} (x-y)^2 \Phi(t, x-y) u_0(y) dx dy + \iint_{\mathbb{R} \times \mathbb{R}} 2y(x-y) \Phi(t, x-y) u_0(y) dy dx$$

$$+ \iint_{\mathbb{R} \times \mathbb{R}} \Phi(t, x-y) y^2 u_0(y) dx dy$$

$$\frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{+\infty} x^2 e^{-\frac{x^2}{4t}} dx = \frac{4t}{\sqrt{\pi}} \int_{-\infty}^{+\infty} \underbrace{t^2 e^{-t^2}}_{-xt \underbrace{te^{-t^2}}_{-\frac{1}{2} \frac{d}{dt} e^{-t^2}}} dt = \frac{4t}{\sqrt{\pi}} \left[-\frac{1}{2} e^{-t^2} \right]_{-R}^R + \frac{1}{2} \int_{-R}^R e^{-t^2} dt$$

$= 2t$

$$= 2t \int_{\mathbb{R}} u_0(y) dy + \int_{\mathbb{R}} y^2 u_0(y) dy$$

$$\frac{d}{dt} \int_{-\infty}^{+\infty} x^2 u(t, x) dx = \int_{-\infty}^{+\infty} x^2 u_t dx = \int_{-\infty}^{+\infty} x^2 u_{xx} dx = - \int_{-\infty}^{+\infty} 2x u_x dx = 2 \int_{-\infty}^{+\infty} u_0(y) dy$$

↑ per parti $[-R, R]$
 $\mathbb{R} \rightarrow +\infty$

③ Risolvere i problemi di Cauchy

$$\begin{cases} u_t = \Delta u \\ u(0, x) = g(x) \end{cases}$$

dove $g_1(x) = \mathbb{1}_{[-1,1]}$, $g_2(x) = \sin x$, $g_3(x) = e^x$

$$\underline{1} \quad u(t, x) = \int_{\mathbb{R}} \Phi(t, x-y) g_1(y) dy = \int_{-1}^1 \frac{1}{\sqrt{4\pi t}} e^{-\frac{(x-y)^2}{4t}} dy$$

pongo $\frac{y-x}{\sqrt{4t}} = s$

$$\text{Erf}(-s) = -\text{Erf}(s)$$

$$= \frac{1}{\sqrt{\pi}} \int_{-\frac{1+x}{\sqrt{4t}}}^{\frac{1-x}{\sqrt{4t}}} e^{-s^2} ds = \frac{1}{2} \left[\text{Erf}\left(\frac{1-x}{\sqrt{4t}}\right) + \text{Erf}\left(\frac{1+x}{\sqrt{4t}}\right) \right]$$



$$\underline{2} \quad u(t, x) = \int_{\mathbb{R}} \frac{1}{\sqrt{4\pi t}} e^{-\frac{(x-y)^2}{4t}} \sin y dy$$

pongo $s = \frac{y-x}{\sqrt{4t}}$ $y = x + \sqrt{4t} s$

$$= \frac{1}{\sqrt{\pi}} \int e^{-s^2} \sin(x + \sqrt{4t} s) ds =$$

$$= \frac{1}{\sqrt{\pi}} \int e^{-s^2} \left[\sin x \cos(\sqrt{4t} s) + \cos x \sin(\sqrt{4t} s) \right] ds$$

$$= \frac{1}{\sqrt{\pi}} \left[\sin x \int e^{-s^2} \cos(\sqrt{4t} s) ds + \cos x \int e^{-s^2} \sin(\sqrt{4t} s) ds \right]$$

$\underbrace{\hspace{10em}}_{v_1(t)}$
 $\underbrace{\hspace{10em}}_{v_2(t) \equiv 0}$

Consideriamo la funzione $t \mapsto \frac{1}{\sqrt{\pi}} \int e^{-s^2} \cos(\sqrt{4t} s) ds = g(t)$

derivato rispetto a t : $g'(t) = \frac{1}{\sqrt{\pi}} \int e^{-s^2} \sin(\sqrt{4t} s) \frac{ds}{\sqrt{4t}} \cdot \frac{d}{dt} (\sqrt{4t} s)$

ora integro per parti in $s \rightarrow$

$$g'(t) = \frac{-1}{\sqrt{\pi}} \int_{-\infty}^{+\infty} \left(\frac{d}{ds} e^{-s^2} \right) \sin(\sqrt{4t} s) \cdot \frac{1}{\sqrt{4t}} ds$$

$$= \frac{-1}{\sqrt{\pi}} \int e^{-s^2} \cos(\sqrt{4t} s) ds = -g(t)$$

inoltre $g(0) = 1 \Rightarrow g(t) = e^{-t}$

$$u(t, x) = e^{-t} \sin x$$

Formalmente, $u(t, x) = \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} e^{-\frac{(x-y)^2}{4t}} e^y dy$
 converge $\forall t > 0$ fissato

$$(x-y)^2 - 4ty = x^2 - 2xy + y^2 - 4ty = x^2 - 2(x+2t)y + y^2 =$$

$$(y - (x+2t))^2 - (x+2t)^2 + x^2 = (y - (x+2t))^2 - 4tx - 4t^2$$

$$-\frac{(x-y)^2 - 4ty}{4t} = -\frac{(y - (x+2t))^2}{4t} + (x+t)$$

$$\Rightarrow u(t,x) = e^{x+t} \frac{1}{\sqrt{4\pi t}} \int e^{-\frac{(y - (x+2t))^2}{4t}} dy = e^{x+t}$$

(4) Se $u(t,x) = v(t)w(x) \Rightarrow u_t - u_{xx} = v'(t)w(x) + v(t)w''(x)$

se $u_t - u_{xx} = 0 \Rightarrow \frac{v'(t)}{v(t)} = -\frac{w''(x)}{w(x)}$

$$\Rightarrow \frac{v'(t)}{v(t)} = \alpha = -\frac{w''(x)}{w}$$

se $\alpha = 0 \Rightarrow v(t) = \text{const}$ e $w'' = 0 \Rightarrow w(x) = ax + b$

e $u(t,x) = a'x + b'$

sol

se $\alpha \neq 0$

$v(t) = A e^{\alpha t}$ e $w(x) = c_1 e^{\sqrt{\alpha}x} + c_2 e^{-\sqrt{\alpha}x}$

se $\alpha > 0 \Rightarrow \alpha = \omega^2$ $u(t,x) = A e^{\omega t} (c_1 e^{\omega x} + c_2 e^{-\omega x})$

se $\alpha < 0 \Rightarrow \alpha = -\omega^2$ $u(t,x) = A e^{-\omega^2 t} (c_1 \cos(\omega x) + c_2 \sin(\omega x))$

(5) se $u(t,x) = \varphi(x-ct)$ resolve $u_t - u_{xx} = 0$

$$\Rightarrow c\varphi' + \varphi'' = 0 \Rightarrow (c\varphi + \varphi')' = 0$$

$$c\varphi + \varphi' = a \Leftrightarrow e^{-cs} (e^{cs} \varphi)' = a \Leftrightarrow e^{cs} \varphi = b + \int_0^1 a e^{ct} dt$$

$$\Rightarrow \varphi(s) = c_1 e^{-cs} + c_2 \frac{a}{c} (e^{ct} - 1)$$

